

Spin Manipulation with an RF Wien-Filter at COSY

PSTP Workshop 2015

Bochum, September 15, 2015 | Sebastian Mey and Ralf Gebel for the JEDI Collaboration |

Forschungszentrum Jülich

Content

EDM Measurements in Magnetic Storage Rings

The Prototype RF ExB-Dipole

Measurements

Summary and Conclusion

Motivation

JEDI Collaboration: First direct measurement of charged light hadrons' permanent **E**lectric **D**ipole **M**oment in storage rings

- simple system with EDM \vec{d} and MDM $\vec{\mu}$ aligned with spin \vec{S}

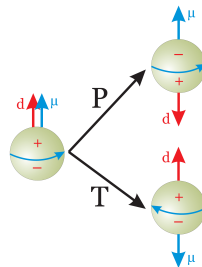
$$\mathcal{H} = -\mu \vec{S} \cdot \vec{B} - d \vec{S} \cdot \vec{E}$$

$$\mathcal{P}(\mathcal{H}) = -\mu \vec{S} \cdot \vec{B} + d \vec{S} \cdot \vec{E}$$

$$\mathcal{T}(\mathcal{H}) = -\mu \vec{S} \cdot \vec{B} + d \vec{S} \cdot \vec{E}$$

⇒ EDMs violate tests **both** parity \mathcal{P} and time reversal \mathcal{T} symmetry

- CPT Theorem: permanent EDMs violate \mathcal{CP} symmetry



source: en.wikipedia.org



Spin Motion in a Magnetic Storage Ring

JEDI Collaboration: First direct measurement of charged light hadrons' permanent EDM in storage rings

- spin motion: $\frac{d\vec{S}}{dt} = \vec{S} \times (\vec{\Omega}_{\text{MDM}} + \vec{\Omega}_{\text{EDM}})$ (Thomas-BMT Equation)

$$\vec{\Omega}_{\text{MDM}} = \frac{q}{m} \left((1 + \gamma \mathbf{G}) \vec{B}_{\perp} + (1 + \mathbf{G}) \vec{B}_{\parallel} - \left(\frac{\gamma}{\gamma+1} + \gamma \mathbf{G} \right) \vec{\beta} \times \vec{E}/c \right)$$

$$\vec{\Omega}_{\text{EDM}} = \frac{q}{m} \frac{\eta}{2} \left(\vec{E}/c + \vec{\beta} \times \vec{B} \right)$$

- MDM: $\vec{\mu} = 2(\mathbf{G} + 1) \frac{q\hbar}{2m} \vec{S}$ with anomalous magnetic moment \mathbf{G}
- EDM: $\vec{d} = \eta \frac{q\hbar}{2mc} \vec{S} \approx 10^{-31} \text{ ecm} \Leftrightarrow \eta \approx 10^{-15}$ for SM light hadrons



Spin Motion in a Magnetic Storage Ring

JEDI Collaboration: First direct measurement of charged light hadrons' permanent EDM in storage rings

- spin motion: $\frac{d\vec{S}}{dt} = \vec{S} \times (\vec{\Omega}_{\text{MDM}} + \vec{\Omega}_{\text{EDM}})$ (Thomas-BMT Equation)

- stationary ring with vertical guiding field \vec{B}_{\perp} and $\vec{B}_{\parallel} = \vec{E} = \vec{0}$

$$\vec{\Omega}_{\text{MDM}} = \frac{q}{m} \left((1 + \gamma G) \vec{B}_{\perp} + (1 + G) \vec{B}_{\parallel} - \left(\frac{\gamma}{\gamma+1} + \gamma G \right) \vec{\beta} \times \vec{E}/c \right)$$

$$\vec{\Omega}_{\text{EDM}} = \frac{q}{m} \frac{\eta}{2} \left(\vec{E}/c + \vec{\beta} \times \vec{B}_{\perp} \right) \text{ couples to motional electric field}$$

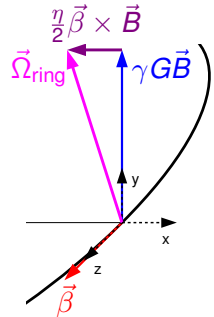
- MDM: $\vec{\mu} = 2(G + 1) \frac{q\hbar}{2m} \vec{S}$ with anomalous magnetic moment G
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Generating an EDM Signal

stationary ring with vertical guiding field \vec{B} and $\vec{B}_{\parallel} = \vec{E} = \vec{0}$

$$\vec{\Omega}_{\text{ring}} = \frac{q}{m} \left((1 + \gamma G) \vec{B} + \frac{\eta}{2} \vec{\beta} \times \vec{B} \right)$$

- spin precession around vertical axis with tune γG
 - tiny EDM tilt of precession axis
 - prepare beam with purely horizontal spins
- ⇒ oscillating vertical spin component, but signal mutch too small to observe

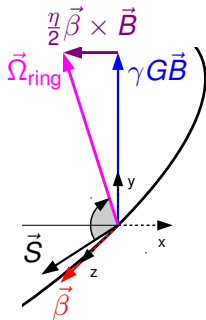


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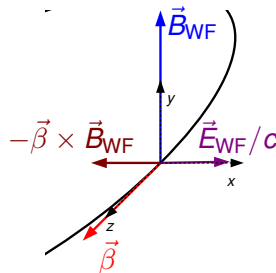


- introduce additional in-plane spin kick in phase with precession
- ⇒ oscillating spins point forward most of the time
- ⇒ **continuous build-up of vertical spin component ⇒ EDM Signal**

Generating an EDM Signal, cont.

- supplement lattice with **local** vertical magnetic field \vec{B}_{WF} oscillating with spin precession
- minimize beam perturbation by adjusting net Lorentz Force to zero

$$\vec{E}_{WF}/c = -\vec{\beta} \times \vec{B}_{WF} \text{ (Wien-Filter condition)}$$



- additional spin rotation in RF Wien-Filter around vertical axis

$$\vec{\Omega}_{MDM} = \frac{q}{m} \left((1 + \gamma G) \vec{B}_{WF} - \left(\frac{\gamma}{\gamma+1} + \gamma G \right) \vec{\beta} \times \vec{E}_{WF}/c \right)$$

$$\vec{\Omega}_{EDM} = \frac{q}{m} \frac{\eta}{2} \left(\vec{E}_{WF}/c + \vec{\beta} \times \vec{B}_{WF} \right) = \vec{0}$$

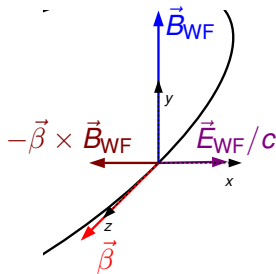
Generating an EDM Signal, cont.

- supplement lattice with **local** vertical magnetic field \vec{B}_{WF} oscillating with spin precession
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- additional spin rotation in RF Wien-Filter around vertical axis

$$\vec{\Omega}_{MDM} = \frac{q}{m} \frac{1+G}{\gamma} \vec{B}_{WF}$$



The RF Wien-Filter itself is EDM transparent, but is capable of generating an EDM signal due to modulation of the spin precession.*

[* W. M. Morse, Y. F. Orlov and Y. K. Semertzidis, Phys. Rev. ST Accel. Beams 16, 114001 (2013)]

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The Prototype RF ExB-Dipole

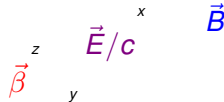
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Prototype RF Wien-Filter with Radial Magnetic Field

$$-\vec{\beta} \times \vec{E}/c$$



- investigate action of RF Wien-Filter fields by direct observation of resulting MDM motion

⇒ use radial magnetic field with vertically prepared spins

⇒ continuous rotation of spin vector during operation

- Lorentz force compensation: $\vec{E}/c = -\vec{\beta} \times \vec{B}$
- spin precession: $\vec{\Omega}_{\text{MDM}} = \frac{1+G}{\gamma} \vec{B}$
- particles sample localized RF field once each turn at orbit angle θ

$$\Rightarrow b(\theta) = \int \hat{B} dz \cos\left(\frac{f_{\text{RF}}}{f_{\text{rev}}} \theta + \phi\right) \sum_{n=-\infty}^{\infty} \delta(\theta - 2\pi n)$$



Resonance Strength of an RF Wien-Filter

- intrinsic resonance strength given by spin rotation per turn, calculate Fourier integral over driving fields along orbit*:

$$\begin{aligned}\epsilon_K &= \frac{f_{\text{spin}}}{f_{\text{rev}}} = \frac{1+G}{2\pi\gamma} \oint \frac{b(\theta)}{B\rho} e^{iK\theta} d\theta \\ &= \frac{1+G}{2\pi\gamma} \frac{\int \hat{B} dz}{B\rho} \sum_{n=-\infty}^{\infty} \cos(2\pi n \frac{f_{\text{RF}}}{f_{\text{rev}}} + \phi) e^{i2\pi K n} \\ &= \frac{1+G}{2 \cdot 2\pi\gamma} \frac{\int \hat{B} dz}{B\rho} \sum_n e^{\pm i\phi} \delta(n - K \mp \frac{f_{\text{RF}}}{f_{\text{rev}}})\end{aligned}$$

- spin tune $\approx \gamma G$, resonance at every sideband with
 $K \stackrel{!}{=} \gamma G = n \pm \frac{f_{\text{RF}}}{f_{\text{rev}}} \Leftrightarrow f_{\text{RF}} = f_{\text{rev}} |n - \gamma G|; n \in \mathbb{Z}$
- d at 970 MeV/c: $f_{\text{rev}} = 750.603 \text{ kHz}; \gamma G = -0.16098$

n	0	1	-1	2	-2
f_{RF}/ kHz	120	629	871	1380	1621

[* S. Y. Lee, 10.1103/PhysRevSTAB.9.074001 (2006)]



The Prototype RF ExB Dipole

coil: 8 windings
length 560 mm

distance 54 mm
length 580 mm

Parameters	RF B dipole
$P_{\text{RMS}} / \text{W}$	90
\hat{I} / A	5
$\int \hat{B}_x dl / \text{Tmm}$	0.175
$f_{\text{RF}} \text{ range} / \text{kHz}$	629 - 1170

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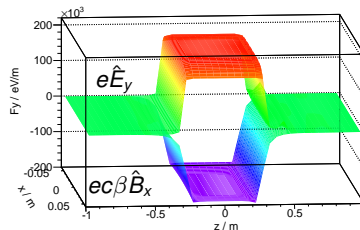


The Prototype RF ExB Dipole

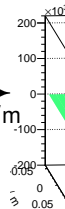
coil: 8 windings

distance 54 mm

length 580 mm



$$\int \hat{F}_y dz = 0 \text{ eV/m}$$



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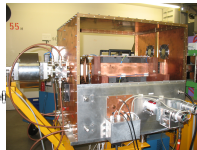
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COSY as Spin Physics R&D Facility

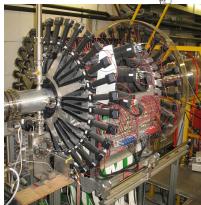


RF solenoid



RF ExB dipole

$\varepsilon_{x,y}$ and $\frac{\Delta p}{p}$ control
electron cooling



fast, continuous
polarimetry



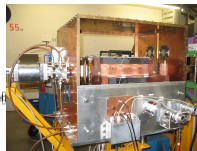
polarized source



COSY as Spin Physics R&D Facility

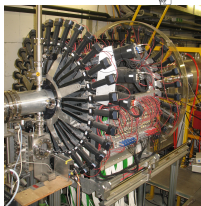


RF solenoid



RF ExB dipole

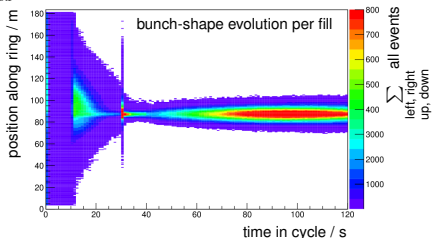
$\varepsilon_{x,y}$ and $\frac{\Delta p}{p}$ control
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RF ExB Setup for Field Compensation

- move betatron sideband onto RF freq. for max. sensitivity

$$q_y \cdot f_{\text{rev}} \stackrel{!}{=} (1 + \gamma G) f_{\text{rev}} = 629 \text{ kHz}$$

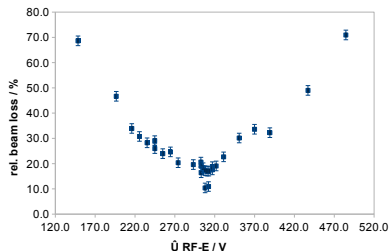
- polarimeter target directly above beam limits acceptance

⇒ exited part of beam is removed

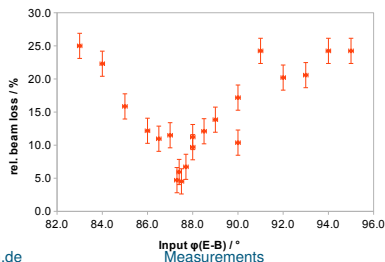
⇒ diagnosis with COSY beam current transformer over $\Delta t = 30 \text{ s}$

- determination of amplitudes and phase corresponding to Lorentz force compensation down to per mille!

Amplitude Scan RF-E at $\hat{I}_{\text{RF-B}} = 2 \text{ A}$



Phase Scan





Beam Response

Analogue signal from one vertical BPM pickup electrode during RF operation exactly **on resonance**

Center $f_{qy} = f_{\text{rev}}(1 + q_y) = 1380 \text{ kHz}$, Span $\Delta f = 10 \text{ kHz}$

RF Wien-Filter:

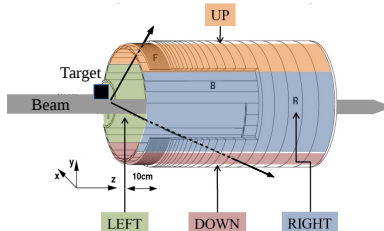
$$\hat{I}_{\text{RF-B}} \approx 740 \text{ mA}; \hat{U}_{\text{RF-E}} \approx 108 \text{ V}$$

RF Sol.:

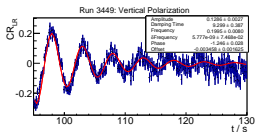
$$\hat{I}_{\text{Sol.}} \approx 780 \text{ mApp}$$

Polarization Measurements

- beam polarization \Leftrightarrow average over all particles' spins
- massive carbon target with slow extraction \Rightarrow long observation time
- polarization signal \Rightarrow rate asymmetries in $^{12}\text{C}(\vec{d}, d)$: $P_y \propto \frac{N_{\text{left}} - N_{\text{right}}}{N_{\text{left}} + N_{\text{right}}}$
- continuous rotation of $\vec{P} \Rightarrow$ oscillation of P_y



Measurement Resonance Strength

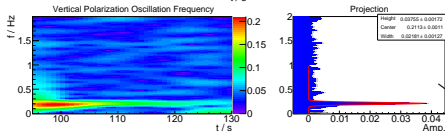


Final $\langle CR_{LR} \rangle = -0.035 \pm 0.005$

Fitted $f_p = (0.1995 \pm 0.0080)$ Hz

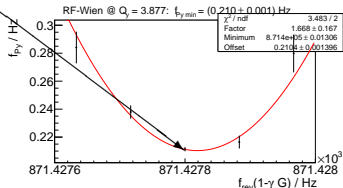
Fitted $\delta f_p = (0.0000 \pm 0.0747)$ Hz

FFT $f_p = (0.2113 \pm 0.0011)$ Hz



- RF Wien-Filter and RF Solenoid both drive continuous rotation of \vec{P}
- find resonance by scan of driving frequency

$$f_{RF} \stackrel{!}{=} f_{rev}(1 - \gamma G)$$

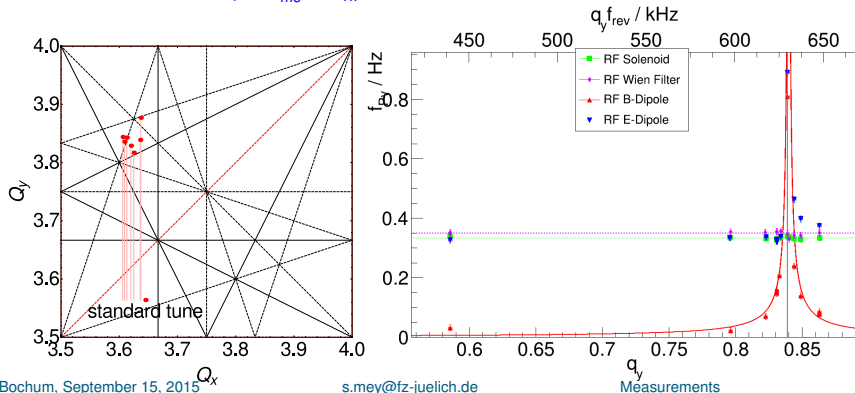


- total spin flip only on resonance \Rightarrow average polarization $\rightarrow 0$
- \Rightarrow minimum of oscillation frequency f_{py}
- measurement of resonance strength

$$\varepsilon = \frac{f_{py \min}}{f_{rev}}$$

Preliminary result of Fixed Frequency Scans

- resonance strength measurements to determine level of field compensation
- RF Solenoid: $f_{Py} = \frac{q}{p} \frac{1+G}{4\pi} \int \hat{B} dI$ RF Wien-Filter: $f_{Py} = \frac{q}{p} \frac{1+G}{4\pi\gamma} \int \hat{B} dI$
- RF B-Dipole: $f_{Py} = \frac{q}{p} \frac{1+\gamma G}{4\pi} \int \hat{B} dI + \text{interference due to beam motion}$
- RF E-Dipole: $f_{Py} = \frac{q}{mc^2} \frac{1/\gamma+1+G}{4\pi} \int \hat{E} dI + \text{interference due to beam motion}$



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Summary

- versatile prototype RF ExB dipole with minimal excitation of coherent beam oscillations has been successfully commissioned
- $P_{\text{RMS}} = 90 \text{ W} \Rightarrow \int \hat{B}_x \text{d}l = 0.175 \text{ T mm}; \int \hat{E}_y \text{d}l = 24.1 \text{ kV}$
Frequency Range 630 kHz - 1060 kHz
- entirely beam-based method for field matching has been worked out and verified
- spin manipulation performance on the same level as with the “proven” RF-Solenoid”system



Outlook

- first attempt of a direct measurement of the deuteron EDM requires a upright, high precision version of an RF Wien-Filter
 - rotatable **stripline** solution scheduled for commissioning at COSY in summer 2016
- ⇒ introduction of the concept and field simulations → J. Slim,
“Towards a High-Accuracy RF Wien Filter for Spin Manipulation at COSY Jülich”

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